



# RELIABLE DYNAMIC ANALYSIS OF AN UNCERTAIN SHEAR BEAM

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# Shear Beams



In structural engineering, shear beams and their dynamic behavior play an important role in

- Modeling,
- Analysis, and
- Design

of various types of structures subjected to a system of dynamic loads such as wind or earthquake excitations.



# Shear Beam Significance



- The dynamic properties of the system are sufficiently simple.
- Therefore, it serves as a convenient representation for illustrating various fundamental features of earthquake response of structures.



# Shear Beam Significance



- It introduces the dynamic analysis of continuous systems with significant shear effects.
- Moreover, the uniform Shear Beam serves as an approximate model for earthquake and wind response of very important types of structures: moderately tall and regular buildings.



# Shear Beam (Cont.)

In current procedures for dynamic analysis of shear beams, the existence of uncertainty in either system's mechanical properties and/or the characteristics of forcing function is not considered.



# Uncertainties



Uncertainties can be attributed to:

- Physical imperfections,
- Modeling inaccuracies, and
- Load-system interaction complexities.

For reliable design, the presence of uncertainties must be included in analysis procedures.



# Interval Uncertainty



- The interval representation is one method to quantify the uncertainty present in a physical system.
- In this representation, the uncertain parameter varies within the interval defined by extreme values.



# Objective



To develop a method for dynamic analysis of a shear beam with properties of structure and load expressed as interval quantities.





# Presentation Outline



- Review of Deterministic Continuous Dynamic Analysis for Shear Beams
- Interval Concepts
- Introduce Interval Continuous dynamic Analysis for Shear Beams
- Numerical Example
- Conclusions



# Shear Beam



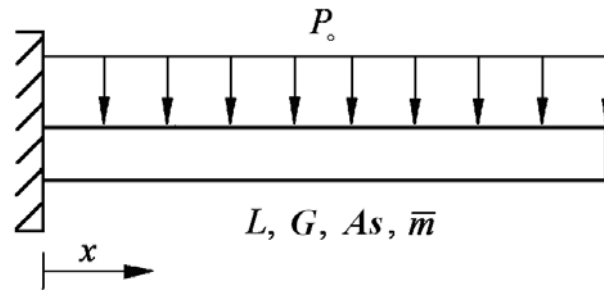
- The shear beam theory was developed by Ukrainian/Russian-born scientist Stephen Timoshenko in the beginning of the 20th century.
- Suitable for describing the behavior of short beams, sandwich composite beams and moderately tall buildings.
- The model takes into account shear deformation and rotational inertia effects.



# Dynamic Deterministic Analysis



Considering a cantilever shear beam subjected to a suddenly applied uniformly distributed load:



The partial differential equation of motion is:

$$\frac{\partial}{\partial x} \left( GA_s \frac{\partial u(x,t)}{\partial x} \right) + P_0 = \bar{m} \frac{\partial^2 u(x,t)}{\partial t^2}$$

# Solution

## Free Vibration



Considering the free vibration and assuming a harmonic function:

$$u(x,t) = \varphi(x)e^{i\omega t}$$

Substitute in the equation of motion, the linear eigenvalue problem:

$$-\frac{\partial}{\partial x} \left( GA_s \frac{d\varphi(x)}{dx} \right) = \bar{m} \omega^2 \varphi(x)$$

Considering:

$$b^2 = \frac{\bar{m} \omega^2}{GA_s}$$



# Solution Free Vibration



Considering a solution:

$$\varphi(x) = A\sin(bx) + B\cos(bx)$$

Applying boundary conditions for simply-supported beam:

$$\varphi(0) = \varphi'(L) = 0$$

The non-trivial solution to the characteristic equation can be obtained.



# Solution

## Free Vibration



Natural circular frequencies (Eigenvalue):

$$\omega_n = (2n - 1) \frac{\pi}{2} \sqrt{\frac{GA_s}{\bar{m}L^2}}$$

Mass-orthonormalized mode shape (Eigenfunction):

$$\phi_n(x) = \sqrt{\frac{2}{\bar{m}L}} \sin\left((2n - 1) \frac{\pi x}{2L}\right)$$

Normalized by:

$$\int_0^L \phi_n(x) \bar{m} \phi_n(x) dx = 1$$



# Solution Forced Vibration



The solution for the forced vibration may be expressed:

$$u(x,t) = \sum_{n=1}^{\infty} y_n(t) \hat{\varphi}_n(x)$$

where,  $y_n$  is the modal coordinate.

Substituting in equation of motion, decoupling by:

- Premultiplying by each mode shape
- Integrating over the domain
- Invoking orthogonality
- Adding modal damping ratio



# Decoupled System



The modal equation of motion is:

$$\ddot{y}_n(t) + 2\zeta_n \omega_n \dot{y}_n(t) + \omega_n^2 y_n(t) = \int_0^L \hat{\phi}_n(x) P_o dx$$

or

$$\ddot{y}_n(t) + 2\zeta_n \omega_n \dot{y}_n(t) + \omega_n^2 y_n(t) = \frac{2P_o}{\sqrt{m}(2n-1)\pi} \left( 1 - \cos\left(\frac{(2n-1)\pi L}{2L}\right) \right)$$

Considering  $\Gamma_n = \frac{2P_o}{\sqrt{m}(2n-1)\pi}$  as the modal participation factor,



# Updated Modal Coordinate



Defining an updated modal coordinate:

$$d_n(t) = \frac{y_n(t)}{\Gamma_n}$$

The updated modal equation:

$$\ddot{d}_n(t) + 2\zeta_n \omega_n \dot{d}_n(t) + \omega_n^2 d_n(t) = 1$$

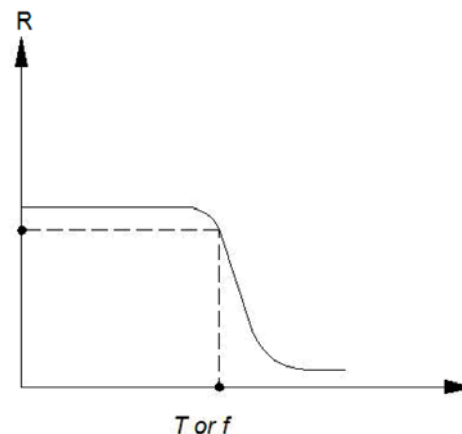


# Maximum Modal Coordinate



The maximum modal coordinate is obtained using response spectrum of each mode for a given  $\omega_n$  and  $\zeta_n$ .

Response Spectrum: Function of maximum dynamic amplification response versus the natural frequencies for an assumed damping ratio.



# Maximum Modal Response



The maximum modal displacement response is the product of:

- Maximum modal coordinate
- Modal participation factor
- Mode shape

$$u_{n,\max} = (d_{n,\max})(\Gamma_n)(\phi_n(x))$$

or:

$$u_{n,\max} = (d_{n,\max}) \frac{8\sqrt{2}P_o}{\bar{m}\sqrt{(2n-1)\pi}} \sin^3\left((2n-1)\frac{\pi x}{4L}\right) \cos\left((2n-1)\frac{\pi x}{4L}\right)$$

# Dynamic Deterministic Analysis

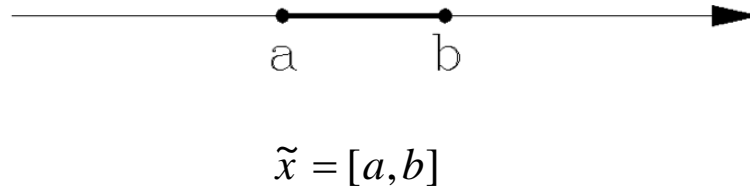


Finally, the total displacement response is obtained using superposition (Square Root of Sum of Squares) of modal maxima:

$$u_{\max} = \sqrt{\sum_{n=1}^N (u_{n,\max})^2}$$

No consideration of uncertainty in the scheme.

# Interval Variable



A real interval is a set of the form:

$$\mathcal{Z} = [z^l, z^u] = \{z \in \mathfrak{R} \mid z^l \leq z \leq z^u\}$$

Archimedes (287-212 B.C.)

$$\left(3\frac{10}{71} < \pi < 3\frac{1}{7}\right)$$

# Interval Dynamic Analysis



Considering a beam with uncertain modulus of elasticity subjected to an uncertain load  $\tilde{G} = [G_l, G_u]$ ,  $\tilde{P}_o = [P_o^l, P_o^u]$ .

The partial differential equation of motion:

$$\frac{\partial}{\partial x} \left( \tilde{G} A_s \frac{\partial u(x,t)}{\partial x} \right) + \tilde{p}(x,t) = m(x) \frac{\partial^2 u}{\partial t^2}$$

# Solution



Natural circular frequencies (Eigenvalue):

$$\tilde{\omega}_n = (2n - 1) \frac{\pi}{2} \sqrt{\frac{\tilde{G}A_s}{mL^2}}$$

Mass-orthonormalized mode shape (Eigenfunction):

$$\phi_n(x) = \sqrt{\frac{2}{mL}} \sin\left((2n - 1) \frac{\pi x}{2L}\right)$$

# Frequencies' Monotonic Behavior



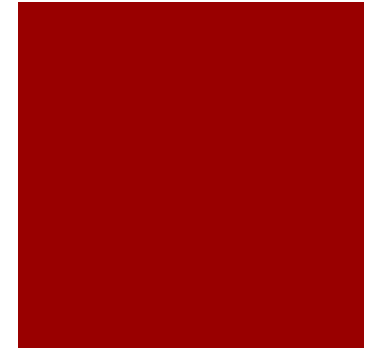
Re-writing interval natural circular frequencies:

$$\omega_n = (2n-1) \frac{\pi}{2} \left[ \sqrt{G^L}, \sqrt{G^U} \right] \cdot \sqrt{\frac{A_s}{mL^2}}$$

This leads to an evident realization of monotonic behavior of natural circular frequencies due to variation in stiffness in continuous dynamic systems.



# Interval Eigenvalue Problem in Discrete Systems



Interval eigenvalue problem using the interval global stiffness matrix:

$$\left(\sum_{i=1}^n ([l_i, u_i])[\bar{K}_i]\right)\{\varphi\} = (\tilde{\omega}^2)[M]\{\varphi\}$$

Rayleigh quotient (ratio of quadratics):

$$R(x) = \frac{x^T Ax}{x^T x}$$

# Bounds on Natural Frequencies



The first eigenvalue – Minimum:

$$\tilde{\lambda}_1 = \min_{x \in R^n} \left[ \frac{x^T \left( \sum_{i=1}^n ([l_i, u_i]) [\bar{K}_i] \right) x}{x^T M x} \right] = \min_{x \in R^n} \left( \sum_{i=1}^n ([l_i, u_i]) \frac{x^T [\bar{K}_i] x}{x^T M x} \right)$$

The next eigenvalues – Maximin Characterization:

$$\tilde{\lambda}_k = \max_{x, z_i=0, i=1, \dots, k-1} \left[ \min_{i=1}^n \frac{x^T \left( \sum_{i=1}^n ([l_i, u_i]) [\bar{K}_i] \right) x}{x^T M x} \right] = \max_{x, z_i=0, i=1, \dots, k-1} \left[ \min_{i=1}^n \sum_{i=1}^n ([l_i, u_i]) \frac{x^T [\bar{K}_i] x}{x^T M x} \right]$$

# Bounding Deterministic Eigenvalue Problems



Solution to interval eigenvalue problem correspond to the maximum and minimum natural frequencies:

$$\left(\sum_{i=1}^n (u_i)[\bar{K}_i]\right)\{\varphi\} = (\omega_{\max}^2)[M]\{\varphi\}$$

$$\left(\sum_{i=1}^n (l_i)[\bar{K}_i]\right)\{\varphi\} = (\omega_{\min}^2)[M]\{\varphi\}$$

Two deterministic problems capable of bounding all natural frequencies of the interval system

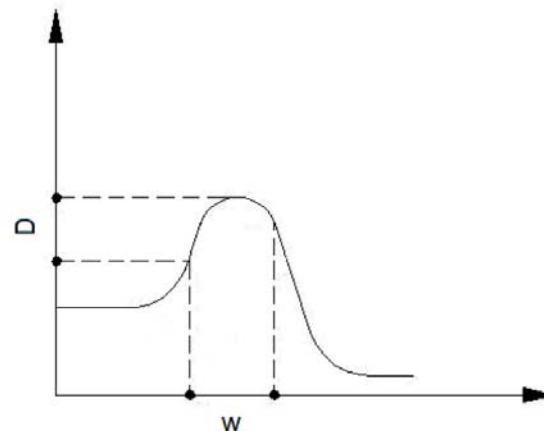
(Modares and Mullen 2004)



# Maximum Modal Coordinate



Having the interval natural frequency, the interval modal coordinate is determined using modal response spectrum as:



The maximum modal coordinate:

$$d_{n,\max} = \max(\hat{d}_n)$$

# Maximum Modal Participation Factor



The interval modal participation factor:

$$\tilde{\Gamma}_n = \frac{2\tilde{P}_o}{\sqrt{m}(2n-1)\pi} \left( 1 - \cos\left(\frac{(2n-1)\pi L}{2L}\right) \right)$$

The maximum modal participation factor:

$$\Gamma_{n,\max} = \frac{2P_o^u}{\sqrt{m}(2n-1)\pi} \left( 1 - \cos\left(\frac{(2n-1)\pi L}{2L}\right) \right)$$

# Maximum Modal Response



The maximum modal displacement response is the product of:

- Maximum modal coordinate
- Modal participation factor
- Mode shape

$$u_{n,\max} = (d_{n,\max})(\Gamma_{n,\max})(\hat{\phi}_n(x))$$

or:

$$u_{n,\max} = (d_{n,\max}) \frac{8\sqrt{2}P_o^u}{\bar{m}\sqrt{(2n-1)\pi}} \sin^3\left((2n-1)\frac{\pi x}{4L}\right) \cos\left((2n-1)\frac{\pi x}{4L}\right)$$

# Deterministic Dynamic Analysis



Finally, the total displacement response is obtained using superposition (Square Root of Sum of Squares) of modal maxima:

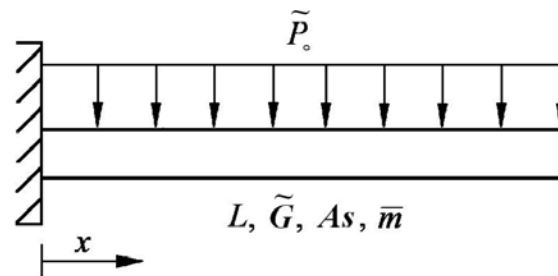
$$u_{\max} = \sqrt{\sum_{n=1}^N (u_{n,\max})^2}$$

Uncertainties are included in the response.

# Numerical Example



Dynamic analysis of a continuous cantilever shear beam with interval uncertainty in the shear modulus and magnitude of uniformly distributed load.



Properties:

$$L = 10 \text{ ft} \quad I = 1530 \text{ in}^4 \quad \bar{m} = 132 \text{ lb/g} \quad \zeta = 1\% \quad \tilde{G} = ([0.9, 1.1])11200 \text{ ksi}$$

$$A_s = 25.23 \text{ in}^2 \quad \tilde{P}_0 = [0.9, 1.1]P_0$$





# Numerical Example



The problem is solved by:

- The present method
- Monte-Carlo simulation  
(Using bounded uniformly distributed random variables in 1000 simulations)



# Results



Bounds on the fundamental natural frequency (first mode)

	Lower Bound <i>Present Method</i>	Lower Bound <i>Monte-Carlo Simulation</i>	Upper Bound <i>Monte-Carlo Simulation</i>	Upper Bound <i>Present Method</i>
$\omega_1$	39.12162	39.12733	43.24335	43.25451

# Numerical Example



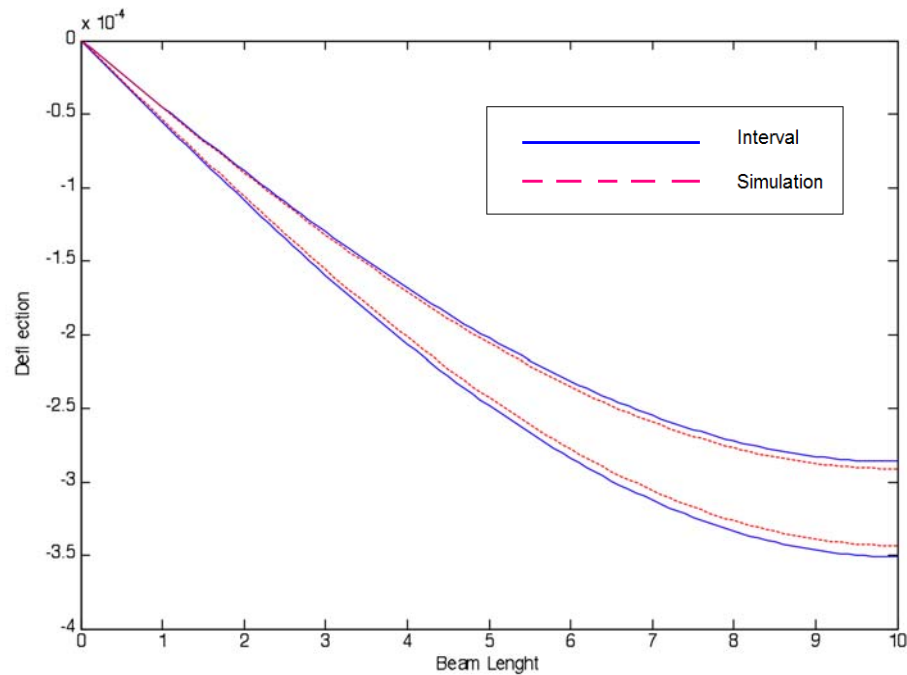
The upper bounds of the maximum displacement response at the beam's free-end

	Upper Bound <i>Monte-Carlo Simulation</i>	Upper Bound <i>Present Method</i>
$\frac{u_{\max}}{P_0}$	0.350714e-3	0.350816e-3

# Results



The first three-mode beam deflection response:



# Sensitivity Analysis

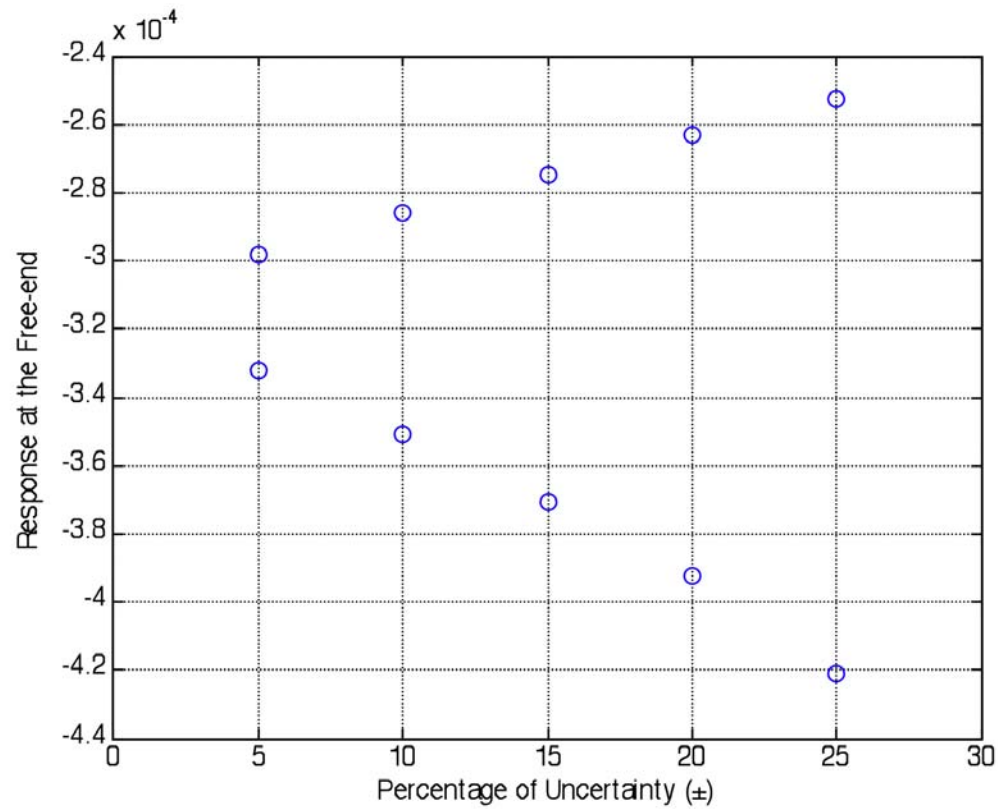


- Performed to investigate the robustness of the proposed method.
- The shear modulus of elasticity is varied as

$$\tilde{G}(n) = 1 + n \times [-0.1, 0.1] \times \text{mid}(\tilde{G})$$

$$n = 1, 2, \dots, 5$$

# Results



# Results



The sensitivity analysis shows that:

- The growth in width of interval of dynamic response is linear.
- In the case of increasing uncertainty, the proposed methodology is robust.



# Conclusions



- A new method for dynamic analysis of shear beam with uncertainty in the mechanical characteristics of the system as well as the properties of the load is developed.
- This computationally efficient method shows that implementation of interval analysis in a continuous dynamic system yields exact and robust results and preserves the problem's physics.





# Conclusions



- The results show that obtaining bounds does not require expensive stochastic procedures such as Monte-Carlo simulations.
- The simplicity of the proposed method makes it attractive to introduce uncertainty in analysis of continuous dynamic systems.



# Questions

